

Densities of Uniform Random Walks on the Plane

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“Probabilists do it on random walks.” – Anonymous

Co-authors: Jon Borwein, Dirk Nuyens, Armin Straub,
and Wadim Zudilin

Outline

- 1 Introduction
 - Numerics 1
- 2 Expectations
 - Experimental maths 1
- 3 Densities
 - Numerics 2
- 4 3 and 4 steps
 - Experimental maths 2

Random walk integrals

Definition: For complex s ,

$$W_n(s) := \int_{[0,1]^n} \left| \sum_{k=1}^n e^{2\pi x_k i} \right|^s d\mathbf{x}$$

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Makes heavy use of experimental mathematics (sophisticated playing).

“Had one not erred, one would have achieved less.” – Euler

Computational challenge

- Dimension reduction:

$$W_n(s) = \int_{[0,1]^{n-1}} \left| 1 + \sum_{k=1}^{n-1} e^{2\pi x_k i} \right|^s d(x_1, \dots, x_{n-1}).$$

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- Tanh-sinh quadrature good for W_3 , after splitting up region so 'bad' points lie on boundaries. Got 175 digits.

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- Binomial expansion:

$$W_n(s) = n^s \sum_{m \geq 0} \frac{(-1)^m}{n^{2m}} \binom{\frac{s}{2}}{m} \underbrace{\int_{[0,1]^n} \left(4 \sum_{i < j} \sin^2(\pi(x_j - x_i)) \right)^m d\mathbf{x}}_{=: I_{n,m}}$$

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- Recursion for $I_{n,m}$ gives accurate evaluations.

Challenge met

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- The recursion satisfies a convolution, which allows rapid computation of odd moments of n steps from $n - 1$ steps (proven for many cases that concern us).

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- The recursion satisfies a convolution, which allows rapid computation of odd moments of n steps from $n - 1$ steps (proven for many cases that concern us).
- We discover these relations because the *computer* allows us to see them.

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Three steps

Theorem (1), Borwein, Nuyens, Straub, W. (2009)

$$W_3(1) = \frac{16\sqrt[3]{4}\pi^2}{\Gamma(\frac{1}{3})^6} + \frac{3\Gamma(\frac{1}{3})^6}{8\sqrt[3]{4}\pi^4},$$
$$W_3(-1) = \frac{3\sqrt[3]{2}}{16\pi^4} \Gamma^6\left(\frac{1}{3}\right).$$

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$$K(x) = \int_0^{\pi/2} \frac{dx}{\sqrt{1-x^2\sin^2(t)}}, \quad E(x) = \int_0^{\pi/2} \sqrt{1-x^2\sin^2(t)} dx.$$

Four steps

Theorem (2), Borwein, Straub, W. (2009)

$$W_4(-1) = \frac{\pi}{4} {}_7F_6 \left(\begin{matrix} \frac{5}{4}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \\ \frac{1}{4}, 1, 1, 1, 1, 1 \end{matrix} \middle| 1 \right).$$

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Theorem (3), Borwein, Straub, W., Zudilin (2010)

$W_4(1) =$

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Where the generalised hypergeometric series is

$${}_pF_q \left(\begin{matrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{matrix} \middle| z \right) = \sum_{n=0}^{\infty} \frac{(a_1)_n \cdots (a_p)_n}{(b_1)_n \cdots (b_q)_n} \frac{z^n}{n!}.$$

Theorem 3 first *identified* by PSLQ.

Cornelius Simon Meijer

Definition (Meijer G-functions)

$$G_{p,q}^{m,n} \left(\begin{matrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{matrix} \middle| x \right) := \frac{1}{2\pi i} \times \int_L \frac{\prod_{j=1}^m \Gamma(b_j - s) \prod_{j=1}^n \Gamma(1 - a_j + s)}{\prod_{j=n+1}^p \Gamma(a_j - s) \prod_{j=m+1}^q \Gamma(1 - b_j + s)} x^s ds.$$

The contour L lies between the poles of $\Gamma(1 - a_i - s)$ and the poles of $\Gamma(b_i + s)$.

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Capture most special functions, very important in CAS, many definite integral evaluations are Meijer G transformations.

$$W_4(s) = \frac{2^s \Gamma(1 + s/2)}{\pi \Gamma(-s/2)} G_{4,4}^{2,2} \left(\begin{matrix} 1, (1 - s)/2, 1, 1 \\ 1/2, -s/2, -s/2, -s/2 \end{matrix} \middle| 1 \right).$$

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- $a(z) := G_{4,4}^{2,2} \left(\begin{matrix} 0, 1, 1, 1 \\ -\frac{1}{2}, \frac{1}{2}, -\frac{1}{2}, -\frac{1}{2} \end{matrix} \middle| z \right)$ not nice; $a(1) = -2\pi W_4(1)$.

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Example of an easy Meijer G identity:

$$\frac{d}{dz} \left(z^{-b_1} G_{4,4}^{2,2} \left(\begin{matrix} a_1, a_2, a_3, a_4 \\ b_1, b_2, b_3, b_4 \end{matrix} \middle| z \right) \right) = \frac{-1}{z^{1+b_1}} G_{4,4}^{2,2} \left(\begin{matrix} a_1, a_2, a_3, a_4 \\ b_1 + 1, b_2, b_3, b_4 \end{matrix} \middle| z \right).$$

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- So probability of returning to the unit disk is

$$\int_0^1 p_n(t)dt = \int_0^\infty J_1(x)J_0^n(x)dx = \left[\frac{-J_0(x)^{n+1}}{n+1} \right]_0^\infty = \frac{1}{n+1}.$$

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- p_n uses radial symmetry. A true 2D density ψ_n satisfies $p_n(x) = 2\pi x\psi_n(x).$

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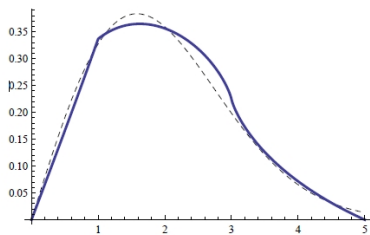
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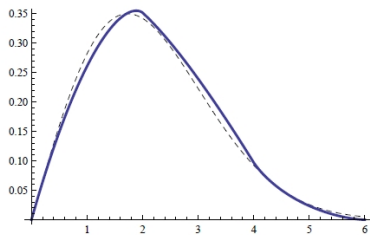
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- $\psi_n(x) \approx \frac{1}{n\pi}e^{-x^2/n}$, like a 2D central limit theorem.

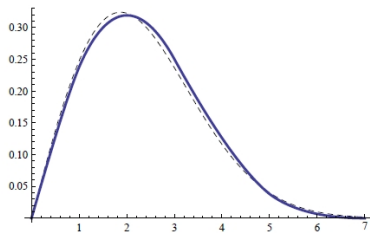
p_n with approximations superimposed.



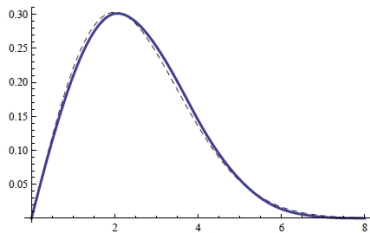
(a) p_5



(b) p_6



(c) p_7



(d) p_8

Karl/Carl Pearson and Eugen Merzbacher

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- Merzbacher (1976): *“The final evaluation must be carried out by numerical computation. . . it has remained difficult to obtain reliable values for p_n when n is moderate. . .”*

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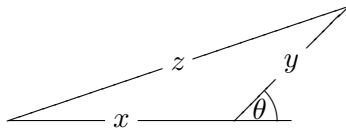
- Asked the original question in 1905.
- *“The graphical construction, however carefully reinvestigated, did not permit of our considering the curve to be anything but a straight line. . .”*
- H. Fettis devotes a whole paper to this.
- A better approximation: $xe^{-\frac{x^2}{n}} \left(\frac{4n^3 - 2n^2 + 4nx^2 - x^4}{2n^4} \right)$ from series for J_0^n .
- Merzbacher (1976): *“The final evaluation must be carried out by numerical computation. . . it has remained difficult to obtain reliable values for p_n when n is moderate. . .”*
- We were able to succeed by looking at moments and analytic structure.

Recursion for p_n

- We condition the distance z of an $(n + m)$ -step walk on x (first n steps), followed by y (remaining m steps).

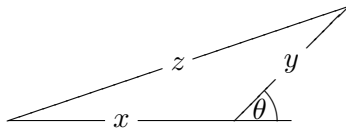
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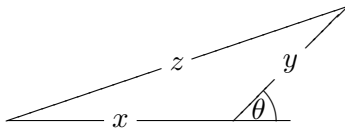


- Therefore

$$W_{n+m}(s) = \frac{1}{\pi} \int_0^n \int_0^m \left(\int_0^\pi z^s d\theta \right) p_n(x) p_m(y) dx dy.$$

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- Change of variable: y as above,

$$W_{n+m}(s) = \int_0^{n+m} z^s \underbrace{\left\{ \int_0^n \int_0^\pi \frac{z}{\pi y} p_n(x) p_m(y) dt dx \right\}}_{p_{n+m}} dz.$$

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- $r = 0 : \psi_n(0) = \psi_{n-1}(1) = \frac{p'_n(0)}{2\pi} = \frac{\operatorname{Res}_{-2} W_n}{2\pi}$.

Outline

- 1 Introduction
 - Numerics 1
- 2 Expectations
 - Experimental maths 1
- 3 Densities
 - Numerics 2
- 4 **3 and 4 steps**
 - Experimental maths 2

Closed form for p_3

- We experimentally found the *modular property*: with the involution $\sigma(x) = \frac{3-x}{1+x}$ (mapping $[0, 1)$ to $(1, 3]$),

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$$p_3(x) = \frac{2\sqrt{3}x}{\pi(3+x^2)} {}_2F_1\left(\frac{1}{3}, \frac{2}{3} \mid \frac{x^2(9-x^2)^2}{(3+x^2)^3}\right),$$

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- It follows that

$$p_3(x) = \frac{2x}{\sqrt{3}\pi} \sum_{k=0}^{\infty} W_3(2k) \left(\frac{x}{3}\right)^{2k}.$$

Poles of W_3 via p_3

- Using this series, we compute (taking care with convergence), for small $\alpha > 0$,

$$\int_0^\alpha p_3(x)x^s dx = \frac{2\alpha^{s+2}}{\sqrt{3}\pi(s+2)} + \frac{2\alpha^{s+4}}{3\sqrt{3}\pi(s+4)} + \dots$$

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- But $W_4(s)$ has double poles. If p_4 admits a Taylor series around the origin, the above argument would produce simple poles. How to reconcile these two facts?

Series for p_4

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- which even I can solve: $f(x) = (a - r \log x)x$, where $a \approx 0.33$.

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- (This method generalises, and 'explains' the shape of p_5 .)

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- DE rigorously produced by *Mellin transforms* of the recursion, *elliptic regularity theorem*, and a *Gosper-like* algorithm.
- More work on modular forms and hypergeometric series led to a closed form for p_4 on $[0, 4]$:

Theorem (4) Borwein, Straub, W., Zudilin (2010)

$$p_4(x) = \frac{2\sqrt{16-x^2}}{\pi^2 x} \operatorname{Re} {}_3F_2 \left(\begin{matrix} \frac{1}{2}, \frac{1}{2}, \frac{1}{2} \\ \frac{5}{6}, \frac{7}{6} \end{matrix} \middle| \frac{(16-x^2)^3}{108x^4} \right).$$

Thank you!

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- Comments?
- Questions?
- Criticisms?